

# Cylinder braiding for quantum linear groups

Tammo tom Dieck

## 1. Introduction and results

The braid group  $ZB_k$  associated to the Coxeter graph  $B_k$  has generators  $t, g_1, \dots, g_{k-1}$  and relations

$$\begin{aligned} tg_1tg_1 &= g_1tg_1t \\ tg_i &= g_it & i > 1 \\ g_i g_j g_i &= g_j g_i g_j & |i - j| = 1 \\ g_i g_j &= g_j g_i & |i - j| > 1. \end{aligned}$$

This group is isomorphic to the group of topological braids with  $k$  strings in the cylinder  $(\mathbb{C} \setminus 0) \times [0, 1]$ .

It is the purpose of this paper to describe a general construction of tensor representations of  $ZB_k$ , based on quantum linear groups.

A tensor representation of  $ZB_k$  is obtained from any *four braid pair*  $(X, F)$  on a module  $V$  over a commutative ring  $\mathfrak{K}$ . By definition, this consists of a Yang-Baxter automorphism (=  $R$ -matrix)  $X: V \otimes V \rightarrow V \otimes V$  which satisfies the braid relation  $(X \otimes 1)(1 \otimes X)(X \otimes 1) = (1 \otimes X)(X \otimes 1)(1 \otimes X)$  on  $V \otimes V \otimes V$  and an automorphism  $F: V \rightarrow V$  such that the *four braid relation*

$$X(F \otimes 1)X(F \otimes 1) = (F \otimes 1)X(F \otimes 1)X$$

holds on  $V \otimes V$ . A representation of  $ZB_k$  on the  $k$ -fold tensor power  $V^{\otimes k}$  is then obtained by the assignment

$$\begin{aligned} t &\mapsto F \otimes 1 \otimes \dots \otimes 1 \\ g_i &\mapsto 1 \otimes \dots \otimes X \otimes \dots \otimes 1 = X_{(i)}, \end{aligned}$$

where the morphism  $X$  in  $X_{(i)}$  acts on the factors  $i$  and  $i + 1$ .

The results of this paper are based on a fundamental four braid pair (1.3) which we now describe. Let  $q \in \mathfrak{K}^\times$  (= the units of  $\mathfrak{K}$ ) and write  $\delta = q - q^{-1}$ . Let

$$P = (p_{ij} \mid 1 \leq i, j \leq n)$$

be a matrix with entries in  $\mathfrak{K}$  such that  $p_{ii} = q$  and  $p_{ij}p_{ji} = 1$  for  $i \neq j$ . Let  $V$  be a free  $\mathfrak{K}$ -module with basis  $v_1, \dots, v_n$ . The assignment

$$(1.1) \quad X(v_i \otimes v_j) = \begin{cases} p_{ij}v_j \otimes v_i & i = j \\ p_{ij}v_j \otimes v_i & i > j \\ p_{ij}v_j \otimes v_i + \delta v_i \otimes v_j & i < j \end{cases}$$

is a Yang-Baxter operator. (Compare [5, p. 171], divide by  $p^{1/2}$  and replace  $q$  by  $p^{-1/2}q$ .) The *standard case*  $p_{ij} = 1$  for  $i \neq j$  is the most interesting. We work with the multi-parameter version, since it makes some of the computations more transparent.

Let  $\beta_j$  ( $1 \leq j \leq n$ ) and  $w$  be elements in  $\mathfrak{K}$  such that

- (1)  $z = \beta_j \beta_{n+1-j}$  is independent of  $j$ ;
- (2)  $\beta_l^2 = w\beta_l + z$  in case  $2l = n + 1$ ;
- (3)  $\prod_j \beta_j \in \mathfrak{K}^\times$ .

Define an automorphism  $F: V \rightarrow V$  by

$$(1.2) \quad F(v_j) = \begin{cases} \beta_j v_{n+1-j} & \text{for } 2j \leq n+1 \\ \beta_j v_{n+1-j} + wv_j & \text{for } 2j > n+1. \end{cases}$$

We use the notation  $j' = n + 1 - j$ . Here is the basic four braid pair; the proof (1.3) will be given in the next section.

**(1.3) Theorem.** *Suppose  $p_{ij}p_{jk} = p_{ij'}p_{j'k}$  whenever  $i \neq j, j'$  and  $k \neq j, j'$ . Then  $(X, F)$ , defined by (1.1) and (1.2), is a four braid pair.*

Henceforth we assume the hypothesis of (1.3).

Associated to the  $R$ -matrix  $X$  above is, by the FRT-construction, a cobraided bialgebra  $A = A(X)$  with *braid form*  $r = r_X: A \otimes A \rightarrow \mathfrak{K}$ . (See [5, VIII 5, VIII 6] or [6, 9.1] for back ground;  $r$  is called universal  $r$ -form in these references.) In our case, the algebra  $A$  is generated by  $T_i^j$ ,  $1 \leq i, j \leq n$ , with relations

$$\begin{aligned} qT_i^m T_i^n &= p_{mn} T_i^n T_i^m \\ qT_i^m T_j^m &= p_{ji} T_j^m T_i^m \\ p_{ji} T_j^m T_i^n &= p_{mn} T_i^n T_j^m \\ \delta T_j^m T_i^n &= p_{mn} T_j^n T_i^m - p_{ij} T_i^m T_j^n \end{aligned}$$

for  $m > n$  and  $i > j$ . The comultiplication  $\mu$  of  $A$  is determined by  $\mu(T_i^j) = \sum_k T_i^k \otimes T_k^j$  and the counit  $\varepsilon$  by  $\varepsilon(T_i^j) = \delta_i^j$ .

It was shown in [2] that any four braid pair  $(X, F)$  induces a so-called cylinder form  $f = f_{(X,F)}: A \rightarrow \mathfrak{K}$ . We recall the notion of a cylinder form on a cobraided bialgebra  $(A, r)$ . We denote the multiplication in the dual algebra  $C^*$  of a coalgebra  $C$  as convolution  $*$ . A *cylinder form* on  $(A, r)$  is a  $\mathfrak{K}$ -linear map  $f: A \rightarrow \mathfrak{K}$  such that  $f$  is convolution invertible and satisfies in the convolution algebra  $(A \otimes A)^*$  the identity

$$(1.4) \quad fm = (f \hat{\otimes} \varepsilon) * r\tau * (\varepsilon \hat{\otimes} f) * r,$$

where  $m: A \otimes A \rightarrow A$  is the multiplication,  $\varepsilon: A \rightarrow \mathfrak{K}$  the counit,  $f \hat{\otimes} g: a \otimes b \mapsto f(a)g(b)$ , and  $\tau(x \otimes y) = y \otimes x$ . In terms of formal notation  $\mu(a) = \sum_{(a)} a_1 \otimes a_2$  and  $(\mu \otimes 1)\mu(a) = \sum_{(a)} a_1 \otimes a_2 \otimes a_3$  for the (iterated) comultiplication, (1.4) assumes the following form

$$f(ab) = \sum_{(a),(b)} f(a_1)r(b_1 \otimes a_2)f(b_2)r(a_3 \otimes b_3)$$

for  $a, b \in A$ .

The quantum linear group  $GL_X(n)$  is a localization of  $A(X)$ , and  $SL_X(n)$  is a quotient of  $A(X)$ . In order to define them, one has to use the quantum determinant  $D$  (see [6, 9.2], [8], [9], [12] for this topic). In our case it is given as a sum over the symmetric group  $S_n$  (compare [3, p. 157])

$$(1.5) \quad D = \sum_{\sigma \in S_n} \alpha(\sigma) T_1^{\sigma(1)} \cdots T_n^{\sigma(n)} = \sum_{\sigma \in S_n} \alpha(\sigma)^{-1} T_n^{\sigma(n)} \cdots T_1^{\sigma(1)}$$

with

$$\alpha(\sigma) = \prod_{\substack{\mu < \nu \\ \sigma(\mu) > \sigma(\nu)}} (-qp_{\sigma(\nu)\sigma(\mu)}).$$

It is not totally obvious that the two expressions of  $D$  define the same element of  $A$ . Since this fact is important for our applications, we indicate a proof as follows. We think of  $A(X)$  as being defined formally over the ring  $\mathbb{Z}[p_{ij}, p_{ij}^{-1}]$ . The algebra has an anti-involution  $\iota$  defined by  $\iota(T_i^j) = T_i^j$ ,  $\iota(q) = q^{-1}$ ,  $\iota(p_{ij}) = p_{ij}^{-1}$ . It sends the first expression for  $D$  into the second one. We have the exterior algebra as a comodule over  $A$ . The exterior algebra is the algebra generated by  $v_1, \dots, v_n$  with relations  $v_i v_i = 0$ ,  $v_i v_j + q^{-1} p_{ij} v_j v_i = 0$  for  $i < j$ . The element  $v_1 v_2 \cdots v_n = v$  generates a one-dimensional subcomodule with coaction  $\mu: v \mapsto D \otimes v$ . We also have an anti-involution  $\iota$  on the exterior algebra which is the identity on the  $v_i$  and defined on the coefficient ring as for  $A$ . The coaction  $\mu$  is compatible with the anti-involutions  $\mu \circ \iota = (\iota \otimes \iota) \circ \mu$ . This compatibility and the fact that  $\iota$  sends  $v$  to a multiple of  $v$  implies  $\iota(D) = D$ .

In the standard case and some other cases, the quantum determinant  $D$  is a group-like central element of  $A$  (see [3, 6.3]). If this is the case, then the algebra  $GL(n) = GL_X(n)$  is defined as the localization of  $A(X)$  with respect to  $D$ , i. e.,  $GL_X(n) = A(X)[y]/(Dy - 1)$ . This is now a Hopf algebra, and the braid form  $r_X$  extends to a braid form on  $GL_X(n)$ ; see [3, 3.1] for a more general result. (In the standard case, these algebras are denoted  $\mathcal{O}(GL_q(n))$  and  $\mathcal{O}(SL_q(n))$  in [6, 9.2].)

We assume  $q^{1/n} \in \mathfrak{K}$  and renormalize our  $R$ -matrix  $X' = q^{-1/n} X$ . We assume that  $D$  is group-like and central. The quotient of  $A(X')$  by the twosided ideal generated by  $D - 1$  is then a Hopf algebra. We denote it by  $SL_X(n) = SL(n)$ . The braid form  $r = r_{X'}$  passes to  $SL(n)$  and yields a braid form.

The braid form  $r = r_{X'}$  on  $A = A(X') = A(X)$  satisfies

$$(1.6) \quad r(x \otimes D) = r(D \otimes x) = \varepsilon(x)$$

for  $x \in A$ . By the axioms of a braid form, we have  $r(xy \otimes D) = r(x \otimes D)r(y \otimes D)$ , since  $D$  is group-like. Therefore it suffices to prove (1.6) for  $x = T_i^j$ . But then it is a restatement of the equality  $\beta = \tau$  in [3, 4.2].

The next proposition collects the formal results about cylinder forms.

**(1.7) Proposition.** *Suppose that the quantum determinant is a group-like central element. Then the cylinder form  $f = f_{(X,F)}$  satisfies  $f(xD) = f(x)f(D)$  for*

$x \in A$ . If  $f(D) \in \mathfrak{K}^\times$ , then the assignment  $f(xD^{-l}) = f(x)f(D)^{-l}$  yields a cylinder form on  $GL(n)$ . If  $f' = f_{(X',F)}$  satisfies  $f'(D) = 1$ , then  $f'$  factors through the quotient map  $A(X') \rightarrow SL(n)$  and induces a cylinder form on  $SL(n)$ .

PROOF. The identity (1.4) together with (1.6) and the fact that  $D$  is group-like yields (in formal notation)

$$\begin{aligned} f(xD) &= \sum_{(x)} f(x_1)r(D \otimes x_2)f(D)r(x_3 \otimes D) \\ &= \sum_{(x)} f(x_1)\varepsilon(x_2)f(D)\varepsilon(x_3) \\ &= f(x)f(D); \end{aligned}$$

the last equality follows from the counit axiom. From  $f(xD) = f(x)f(D)$  we see that  $xD^{-s} \mapsto f(x)f(D)^{-s}$  is a well-defined map on  $GL(n)$ . By using the same property, one verifies that (1.4) still holds for the extended form on  $GL(n)$ . Finally, it is clear that  $f$  vanishes on the ideal generated by  $D - 1$  if  $f(D) = 1$ .  $\square$

In order to make use of (1.7) we have to know the value  $f(D)$ . Theorem (1.8) will be obtained from a general analysis of tensor representations in section 3.

**(1.8) Theorem.** *The cylinder form  $f_{(X,F)}$  assumes the following value on the quantum determinant*

$$f(D) = (-q)^{-n(n-1)/2} \prod_{i>j} p_{i'j} \prod_j \beta_j.$$

In the standard case  $p_{ij} = 1$  for  $i \neq j$ , we have

$$\prod_{i>j} p_{i'j} = q^l \quad \text{for } n = 2l, n = 2l + 1.$$

If one works with  $X'$  in place of  $X$ , then this value has to be multiplied by  $q^{1-n}$ , since the definition of  $f$  via  $t_n$  (see section 3) involves  $n(n-1)$   $R$ -matrices  $X$ .

Finally, we recall from [2] that a cylinder form  $f$  on a cobraided bialgebra  $(A, r)$  induces a cylinder braiding on the category of left  $A$ -comodules. Let  $\mu_M: M \rightarrow A \otimes M$ ,  $x \mapsto \sum x^1 \otimes x^2$  be a comodule structure on  $M$  in formal notation. Define  $t_M: M \rightarrow M$  by  $x \mapsto \sum f(x^1)x^2$ . Let  $z_{M,N}: M \otimes N \rightarrow N \otimes M$  denote the braiding morphisms induced by the braid form  $r$ . Then for any two left  $A$ -comodules  $M$  and  $N$  the following identities hold

$$t_{M \otimes N} = z_{N,M}(t_N \otimes 1)z_{M,N}(t_M \otimes 1) = (t_M \otimes 1)z_{N,M}(t_N \otimes 1)z_{M,N}.$$

Isomorphisms  $t_M: M \rightarrow M$  which satisfy these identities are called a *cylinder braiding* on the category of left  $A$ -comodules. In particular, we obtain a four braid pair  $(z_{M,M}, t_M)$  for each comodule  $M$ .

We remark that the algebra of cylinder braidings can be used to construct invariant of coloured ribbon tangles in the cylinder along the lines of Turaev [13].

## 2. The four braid pair

This section is devoted to the proof of (1.3). Consider the involutions  $\sigma(i, j) = (n + 1 - i, j)$  and  $\tau(i, j) = (j, i)$  on the set of indices  $J = \{(i, j) \mid 1 \leq i, j \leq n\}$ . Since  $\sigma\tau\sigma\tau = \tau\sigma\tau\sigma$ , they formally generate the dihedral group  $D_8$  of order 8. We decompose  $J$  into the orbits under the  $D_8$ -action. Recall the notation  $n+1-i = i'$ .

Let  $n = 2k$ . Then we have the orbit  $(i, i), (i', i), (i, i'), (i', i')$  of length 4. This is the orbit of  $(i, j)$  if  $i = j$  or  $i = j'$ . If  $i \neq j, j'$ , then the orbit of  $(i, j)$  has length 8.

Let  $n = 2\ell - 1$ . Then we have the fixed point  $(\ell, \ell)$ . There is another orbit type of length 4, namely  $(\ell, j), (j, \ell), (j', \ell), (\ell, j')$  for  $j \neq \ell$ . In addition, we have the orbits of  $(i, i)$ ,  $i \neq \ell$ , of length 4; and for  $i \neq j, j'$  and  $i \neq \ell \neq j$  the orbit of  $(i, j)$  of length 8.

For each  $D_8$ -orbit  $x \in J/D_8$  the subspace of  $V \otimes V$  spanned by  $\{v_i \otimes v_j \mid (i, j) \in x\}$  is invariant under  $X$  and  $F \otimes 1 = Y$ .

Therefore it suffices to verify the four braid relation on these subspaces. The structure of the matrices involved depends only on the  $D_8$ -isomorphism type of the orbit. When we display matrices we use the antilex(icographical) order:  $(i, j) < (a, b)$  if and only if  $j < b$  or  $j = b$  and  $i > a$ . This ordering is applied to the basis  $\{v_i \otimes v_j \mid (i, j) \in J\}$  of  $V \otimes V$ .

The fixed point orbit is trivial. Therefore we have to consider three cases.

Case (I): The orbit  $(i, i), (j, i), (i, j), (j, j)$  with  $i < j = i'$ . The matrices are

$$Y = \begin{pmatrix} 0 & \beta_j & & \\ \beta_i & w & & \\ & & 0 & \beta_j \\ & & \beta_i & w \end{pmatrix} \quad X = \begin{pmatrix} q & & & \\ & 0 & p_{ij} & \\ & p_{ji} & \delta & \\ & & & q \end{pmatrix}.$$

One computes  $YXYX$  to be the matrix of  $2 \times 2$ -blocks

$$\begin{pmatrix} 0 & qp_{ij}\beta_j F_{ij} \\ q\beta_i p_{ji} F_{ij} & (q^2 - 1)F_{ij}^2 + wF_{ij} \end{pmatrix} \quad \text{with} \quad F_{ij} = \begin{pmatrix} 0 & \beta_j \\ \beta_i & w \end{pmatrix}.$$

One obtains  $XYXY$  by transposition of  $YXYX$  and the changes  $p_{ab} \mapsto p_{ba}$ ,  $\beta_a \mapsto \beta_{a'}$ . This shows the four braid relation this case.

Case (II). The orbit  $(\ell, i), (i, \ell), (i', \ell), (\ell, i')$  for  $i < \ell$  if  $n = 2\ell - 1$ . The matrices are ( $j = i'$ )

$$Y = \begin{pmatrix} \beta_\ell & & & \\ & 0 & \beta_j & \\ & \beta_i & w & \\ & & & \beta_\ell \end{pmatrix} \quad X = \begin{pmatrix} 0 & p_{i\ell} & & \\ p_{\ell i} & \delta & & \\ & & 0 & p_{\ell j} \\ & & p_{j\ell} & \delta \end{pmatrix}.$$

Recall that  $\beta_\ell^2 = \beta_\ell w + \beta_i \beta_j$ . (Strangely enough,  $Y$  is now an  $R$ -matrix and  $X$  has almost the form of the  $Y$  in case (I).) The necessity of the condition

$p_{i\ell}p_{\ell i} = p_{j\ell}p_{\ell j}$  results from the comparison of the  $2 \times 2$ -blocks on the codiagonal of  $XYXY = YXYX$ . If we invoke this condition and use  $\beta_\ell^2 = \beta_\ell w + z$ , we compute  $XYXY$  to be the matrix of  $2 \times 2$ -blocks

$$\begin{pmatrix} 0 & \beta_j \beta_\ell p_{\ell j} P_{j\ell}^{i\ell} \\ \beta_i \beta_\ell p_{j\ell} P_{\ell i}^{\ell j} & \beta_\ell w (P_{j\ell}^{\ell j})^2 + \delta \beta_i \beta_j P_{j\ell}^{\ell j} \end{pmatrix} \quad \text{with} \quad P_{ab}^{cd} = \begin{pmatrix} 0 & p_{cd} \\ p_{ab} & \delta \end{pmatrix}.$$

Again, this matrix does not change under transposition and the changes as in case (I).

Case (III). The orbit of length eight for  $a < b \leq [\frac{n}{2}]$  is  $(b, a)$ ,  $(b', a)$ ,  $(a, b)$ ,  $(a', b)$ ,  $(a, b')$ ,  $(a', b')$ ,  $(b, a')$ ,  $(b', a')$ . For simplicity of notation we use  $n = 4$ ,  $(a, b) = (1, 2)$ . The matrices are in  $2 \times 2$ -block form

$$Y = \begin{pmatrix} Y_{23} & & & \\ & Y_{14} & & \\ & & Y_{14} & \\ & & & Y_{23} \end{pmatrix} \quad Y_{ab} = \begin{pmatrix} 0 & \beta_b \\ \beta_a & w \end{pmatrix}$$

$$X = \begin{pmatrix} & p_{12} & & \\ p_{21} & \delta & & \\ & & p_{13} & \\ p_{31} & & \delta & p_{24} \\ & & & p_{34} \\ & p_{42} & & \delta \\ & & p_{43} & \delta \end{pmatrix}.$$

The necessity of the condition  $p_{ai}p_{ib} = p_{ai'}p_{i'b}$  results from the comparison of the  $2 \times 2$ -blocks on the codiagonal of  $YXYX$  and  $XYXY$ . If we invoke this condition, then  $YXYX$  turns out to be the following matrix of  $2 \times 2$ -blocks.

			$p_{13}p_{34}\beta_4 Y_{23}$
		$p_{24}p_{43} Y_{14}$	$\delta p_{24}\beta_3 Y_{24}$
	$p_{34}p_{42}\beta_2 Y_{14}$	$w Y_{14}$	$\delta p_{34}(zI + w Y_{24})$
$p_{43}p_{31}\beta_1 Y_{23}$	$\delta p_{42}\beta_2 Y_{13}$	$\delta p_{43}(zI + w Y_{13})$	$\delta^2 zI + (w + \delta^2 w) Y_{23}$

Here  $I$  is a unit matrix. Again, transposition and the changes of case (I) do not change the matrix.

### 3. The cylinder twist in the tensor representation

Let  $E_{ij}: V \rightarrow V$  be the linear map defined by  $E_{ij}(v_r) = \delta_{ir}v_j$ . We use the same symbol  $E_{ij}$  for its matrix, it has 1 in position  $(j, i)$  and 0 in the other positions. If  $M: V \rightarrow V$  is any linear map with matrix  $M = (m_{ab})$ , then

$$(3.1) \quad E_{ji}ME_{lk} = m_{jk}E_{li}.$$

We also use

$$E(r) = \sum_{j=1}^{r-1} E_{jj}, \quad 1 \leq r \leq n.$$

This is the zero matrix for  $r = 1$ .

In order to organize the following computations, we need a suitable notation for matrices in tensor products. Suppose  $U$  and  $W$  are  $\mathfrak{K}$ -modules with basis  $u_1, \dots, u_r$  and  $w_1, \dots, w_s$ . We call

$$u_1 \otimes w_1, \dots, u_r \otimes w_1; u_1 \otimes w_2, \dots, u_r \otimes w_2; \dots u_r \otimes w_s$$

the *antilex*(icographical) basis of  $U \otimes W$ . This ordering of a basis is an associative process. In particular,  $V(k) := V^{\otimes k}$  has for any  $k \in \mathbb{N}$  a well-defined antilex basis. In the sequel, we always use the antilex basis, unless the contrary is explicitly stated.

Suppose  $f: U \rightarrow U$  has matrix  $A$ . Then  $f \otimes 1_W: U \otimes W \rightarrow U \otimes W$  has the block diagonal matrix  $\text{Dia}(A, \dots, A)$  with  $\text{rank}(W)$  diagonal terms  $A$ .

Suppose  $g: W \rightarrow W$  has matrix  $B = (b_{ij})$ . Then  $1_U \otimes g: U \otimes W \rightarrow U \otimes W$  has the block matrix  $B$ , where  $b_{ij}$  now stands for the diagonal block  $b_{ij}I(u)$ , with  $I(u)$  the unit matrix of size  $u = \text{rank}(U)$ .

The matrix  $X$  in (1.1) is in this notation the  $n \times n$ -block matrix with blocks of size  $n$

$$(3.2) \quad X = \begin{pmatrix} qE_{11} + \delta E(1) & p_{12}E_{12} & p_{13}E_{13} & \cdots & p_{1n}E_{1n} \\ p_{21}E_{21} & qE_{22} + \delta E(2) & p_{23}E_{23} & \cdots & p_{2n}E_{2n} \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ p_{n1}E_{n1} & p_{n2}E_{n2} & p_{n3}E_{n3} & \cdots & qE_{nn} + \delta E(n) \end{pmatrix}.$$

By the remarks above, the matrix

$$X_k = 1_{V(k-2)} \otimes X$$

on  $V^{\otimes k}$  has a similar structure, but each entry is now a diagonal block of size  $n^{k-2}$ .

We still use the notation as in (3.2). Then  $E_{ij}$  is a block-matrix of size  $n$ , and its entries are blocks of size  $n^{k-2}$ . In this block notation (3.1) still holds.

Basis objects for our inductive computations are bottom-right triangular block-matrices with blocks of size  $m$ , or BRTB( $m$ )-matrices for short. These are, by definition,  $n \times n$ -matrices  $T = (T_{ij})$  with:

$$(1) \quad T_{ij} = 0 \quad \text{for } i + j \leq n.$$

- (2)  $T_{ij}$  is a block of size  $m$ .
- (3) The blocks  $T_{i,n+1-i}$  on the codiagonal are diagonal matrices.

The next proposition provides us with the basic inductive structure of BRTB-matrices with regard to the matrix  $X$ .

**(3.3) Proposition.** *Let  $T = (T_{ij})$  be a BRTB( $n^{k-2}$ )-matrix. Then*

$$\tilde{T} = (\tilde{T}_{ab}) = X_k(T \otimes 1)X_k$$

*is a BRTB( $n^{k-1}$ )-matrix.*

PROOF. We use the remarks above about the structure of matrices in tensor products, the structure of  $X_k$ , and the relation (3.1). Denote the  $n \times n$ -matrix  $X$  in (3.2) by  $(e_{ab})$ . Then, by matrix multiplication,

$$(3.4) \quad \tilde{T}_{ab} = \sum_{j=1}^n e_{aj}T e_{jb}.$$

We use (3.1) with  $e_{aj} = p_{aj}E_{aj}$ , for  $a \neq j$ , and  $e_{aa} = qE_{aa} + \delta \sum_{j=1}^{a-1} E_{jj}$ . Suppose  $a + b \leq n$ . Then, by assumption,  $T_{ab} = 0$ , and (3.1) gives  $\tilde{T}_{ab} = 0$ .

It remains to show that  $\tilde{T}_{ab}$  for  $a+b = n+1$  is a diagonal matrix. From (3.1) and (3.4) we obtain  $\tilde{T}_{ab} = \sum_j p_{aj}p_{jb}T_{ab}E_{jj}$ , since  $\delta E(a)T e_{ab} = 0$  and  $e_{ab}T \delta E(b) = 0$  by the BRTB( $n^{k-2}$ )-property of  $T$ . This is the diagonal block matrix

$$(3.5) \quad \text{Dia}(p_{a1}p_{1b}T_{ab}, p_{a2}p_{2b}T_{ab}, \dots, p_{an}p_{nb}T_{ab}).$$

We apply this result to the tensor representation obtained from the four braid pair (1.3). We set  $t(1) = F$  and define inductively

$$t(k) = X_k(t(k-1) \otimes 1)X_k.$$

We can apply (3.3) to the matrices  $t(k)$ . We determine the diagonal blocks on the codiagonal of  $t(k)$ . Since we are working in  $V^{\otimes k}$ , we use a multi-index notation for the basis and other terms. A multi-index is a function  $i: \{1, \dots, k\} \rightarrow \{1, \dots, n\}$ . Let  $A(k, n)$  denote the set of these functions. We write

$$v_i = v_{i(1)} \otimes \dots \otimes v_{i(k)} \quad \text{for } i \in A(k, n).$$

The diagonal positions of the diagonal block  $t(k)_{a,n+1-a}$  of  $t(k)$  are indexed by multi-indices  $i$  with  $i(k) = a$ ; here  $i$  denotes the columns of  $t(k)$ . Let  $d(i)$  denote this diagonal term. From (3.5) we obtain by induction on  $k$ :

$$(3.6) \quad d(i) = \beta_{i(k)} \prod_{j=1}^{k-1} p_{i(k)i(j)} p_{i(j)i(k)}.$$

The *cylinder twist*  $t_k: V^{\otimes k} \rightarrow V^{\otimes k}$  is, by definition, the automorphism defined inductively by

$$t_1 = t(1) \quad \text{and} \quad t_k = (t_{k-1} \otimes 1)t(k).$$



**(3.7) Theorem.** *The matrix of  $t_k$  is a BRTB(1)-matrix. The codiagonal element  $\ell(i)$  in column  $i \in A(k, n)$  is*

$$\ell(i) = \prod_{a < b} p_{i(b)i(a)} p_{i(a)i(b)'} \cdot \prod_{a=1}^k \beta_{i(a)}.$$

PROOF. The blocks of  $t_k$  are  $t_{k-1}t(k)_{ab}$ . This immediately gives that  $t_k$  is a BRTB(1)-matrix. The value of the codiagonal element is obtained inductively from (3.6).  $\square$

We mention the standard special case  $p_{ij} = 1$  for  $i \neq j$ . Let  $\text{Or}(n)$  be the orbit set of  $\{1, \dots, n\}$  under the involution  $i \mapsto n + 1 - i$ . Let  $s(x)$  be the order of the isotropy group of  $x \in \text{Or}(n)$ . In that case we have

$$(3.8) \quad \ell(i) = q^{\lambda(i)} \prod_{a=1}^k \beta_{i(a)} \quad \text{with} \quad \lambda(i) = \sum_{x \in \text{Or}(n)} s(x) \binom{|i^{-1}(x)|}{2}.$$

It may be helpful to see a specific example for these computations. We consider the case  $n = 3$  and  $p_{ij} = 1$  for  $i \neq j$ . Then we obtain the matrix

$$t(2) = \begin{array}{c|c|c} & & \begin{array}{c} q\beta_3 \\ \beta_3 \\ q\beta_3 \end{array} \\ \hline & \begin{array}{c} \beta_2 \\ q^2\beta_2 \\ \beta_2 \end{array} & \begin{array}{c} \delta\beta_3 \\ \delta\beta_2 \end{array} \\ \hline \begin{array}{c} q\beta_1 \\ \beta_1 \\ q\beta_1 \end{array} & \begin{array}{c} \delta\beta_1 \\ \delta\beta_2 \end{array} & \begin{array}{c} w \\ w + \delta^2\beta_2 \\ q\delta\beta_1 \end{array} \quad \begin{array}{c} q\delta\beta_3 \\ q^2w \end{array} \end{array}$$

and finally the matrix for  $t_2 = (t_1 \otimes 1)t(2)$  in the antilex basis.

			$q\beta_3^2$
		$q\beta_1\beta_3$	$\beta_2\beta_3$
	$\beta_2\beta_3$		$q\beta_3w$
	$q^2\beta_2^2$		$\delta\beta_2\beta_3$
	$\beta_1\beta_2$	$w\beta_2$	$\delta\beta_2^2$
$q\beta_1\beta_3$		$q\delta\beta_1\beta_3$	$q^2w\beta_3$
$\beta_1\beta_2$	$\delta\beta_1\beta_2$	$\delta\beta_2^2$	$w\beta_2 + \delta^2\beta_2^2$
$q\beta_1^2$	$q\beta_1w$	$q^2w\beta_1$	$q\delta\beta_1\beta_3 + q^2w^2$

The space  $V \otimes V$  decomposes into the second symmetric and exterior power  $V \otimes V = S^2V \oplus \Lambda^2V$ . A basis of  $S^2V$  is  $v_{22}; u_{12}, u_{23}; v_{11}, u_{13}, v_{33}$  with  $v_{ij} = v_i \otimes v_j$

and  $u_{ij} = v_{ij} + q^{-1}v_{ji}$ . The subspaces  $S^2V$  and  $\Lambda^2V$  are  $t_2$ -invariant. The matrix of  $t_2$  on  $S^2V$  in the basis above is a block diagonal matrix with blocks

$$(q^2\beta_2^2), \quad \begin{pmatrix} & q\beta_2\beta_3 \\ q\beta_1\beta_2 & q\delta\beta_2^2 + w\beta_2 \end{pmatrix}, \quad \begin{pmatrix} & q\beta_3^2 \\ q\beta_1^2 & (1+q^2)w\beta_1 & q^2w\beta_3 \\ & q^2w^2 + q\delta\beta_1\beta_3 \end{pmatrix}.$$

The roots of the characteristic polynomial are  $q^2\beta_2^2$ ,  $-\beta_1\beta_3$  and  $q^2(\beta_2 - w)^2$ , with multiplicities 3, 2, and 1, respectively. Thus the matrix satisfies a relation of degree three.

The space  $\Lambda^2V$  has the basis  $w_{ij} = v_{ij} - qv_{ji}$ ,  $i < j$ . The restriction of  $t_2$  to  $\Lambda^2V$  has with respect to the basis  $w_{12}, w_{13}, w_{23}$  the matrix

$$\begin{pmatrix} & & -q^{-1}\beta_2\beta_3 \\ -q^{-1}\beta_1\beta_2 & -\beta_1\beta_3 & \\ & -\beta_1\beta_3 + q^{-2}\beta_2^2 & \end{pmatrix}.$$

The characteristic polynomial has the roots  $-\beta_1\beta_3, q^{-2}\beta_2^2$  with multiplicities 2, 1.

*Proof* of (1.8). We apply the preceding results in order to compute the value of the cylinder form on the quantum determinant  $D$ . For this purpose we have to recall from [2] the definition of the cylinder form  $f: A \rightarrow \mathfrak{K}$ . Suppose  $t_k(v_i) = \sum_j F_i^j v_j$ , where  $i$  and  $j$  are multi-indices of length  $k$ . The cylinder form is then defined by  $f(T_i^j) = F_i^j$ .

We see from Theorem (3.7) that  $f(T_n^{\sigma(n)} \cdots T_1^{\sigma(1)}) = 0$  if  $(\sigma(n), \dots, \sigma(1)) \neq (1, \dots, n)$ . This claim is proved as follows: We have to consider the column  $(n, \dots, 1)$  of the matrix (3.7) for  $t_n$ . Among the  $(\sigma(n), \dots, \sigma(1))$  the element  $(1, \dots, n)$  is the largest one in the antilex order. The claim thus follows from the BRTB-structure of the matrix. Thus only the term  $f(T_n^1 T_{n-1}^2 \cdots T_1^n)$  contributes to  $f(D)$ . The value for  $f(D)$ , as stated in (1.8), is now a direct consequence of (3.7). Note that there is some cancellation of  $p_{ij}$ -factors due to the hypothesis of (1.3).  $\square$

By dualization and linear algebra (see e. g. [10], [12]) one can obtain from the results of this paper cylinder braidings on suitable categories of modules over quantum enveloping algebras (say integrable modules over  $U_q(\mathfrak{sl}_n)$  in the sense of [7]).

## References

1. tom Dieck, T.: On tensor representations of knot algebras. Manuscripta math. 93, 163 – 176 (1997).
2. tom Dieck, T., and R. Häring-Oldenburg: Quantum groups and cylinder braiding. Forum math. (1998).

3. Hayashi, T.: Quantum groups and quantum determinants. *J. of Algebra* 152, 146-165 (1992).
4. Jantzen, J.C.: Lectures on quantum groups. Graduate studies in Math. Vol. 6. Amer. Math. Soc. 1996.
5. Kassel, Ch.: Quantum groups. New York, Springer 1995.
6. Klimyk, A., and K. Schmüdgen: Quantum groups and their representations. New York, Springer 1997.
7. Lusztig, G.: Introduction to quantum groups. Basel, Birkhäuser 1993.
8. Manin, Yu. I.: Quantum groups and non-commutative geometry. CRM, Montreal 1988.
9. Parshall, B., and Jiang-pan Wang: Quantum linear groups. *Mem. Amer. Math. Soc.* 439 (1991).
10. Tanisaki, T.: Finite dimensional representations of quantum groups. *Osaka J. Math.* 28, 37 – 53 (1991).
11. Takeuchi, M.: Matric bialgebras and quantum linear groups. *Israel J. Math.* 72, 232 – 251 (1990).
12. Takeuchi, M.: Some topics on  $GL_q(n)$ . *J. of Algebra* 147, 379 – 410 (1992).
13. Turaev, V.G.: Quantum invariants of knots and 3-manifolds. W. de Gruyter, Berlin 1994.

Version of February 9, 2009

Tammo tom Dieck  
Mathematisches Institut  
Bunsenstr. 3/5  
D – 37073 Göttingen  
tammo@uni-math.gwdg.de